



Groupement
de recherche

ARCHI-META

Architected Metamaterials

Topological lattices for efficient light transport

T. Ozawa, H. M. Price, A. Amo et al., Rev. Mod. Phys. 91, 15006 (2019)

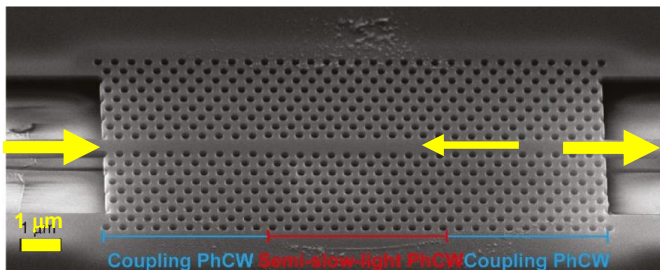
Journée de lancement

27 novembre 2023

Alberto Amo

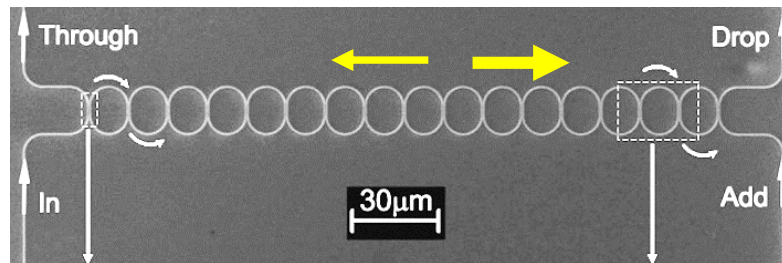
Laboratoire PhLAM – CNRS – Université de Lille

Photonic crystal waveguides



Z. Cheng et al., Nanophotonics 9, 2377 (2020)

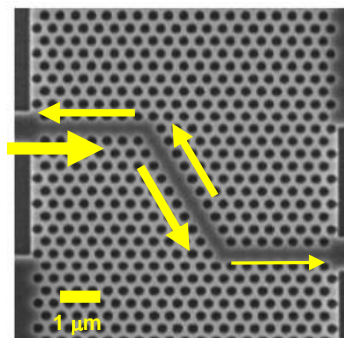
Coupled resonator waveguides (delay lines)



F. Xia et al., Appl. Phys. Lett. 89, 041122 (2006)

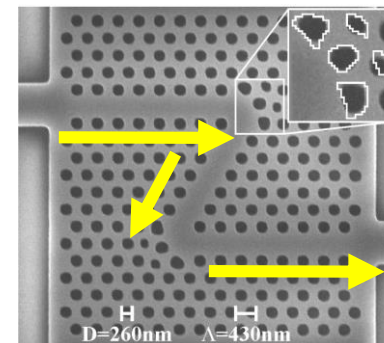
Problems

- ➔ Back scattering at imperfections
- ➔ Very strong bending losses



L. H. Frandsen et al., Opt. Exp. 12, 5916 (2004)

Inverse design using variational algorithms, deep learning etc.



P. I. Borel et al., Opt. Exp. 12, 1996 (2004)

Highly demanding

Topological invariants: a closed surface

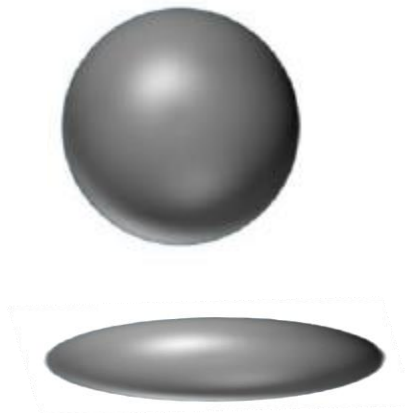
Properties that remain unaffected under smooth distortions

Genus

$$g = 1 - \frac{1}{4\pi} \int_S K(\mathbf{r}) dA = 0, 1, 2, \dots$$

- ➔ $K(\mathbf{r})$ local curvature
- ➔ Number of holes
- ➔ Global property
- ➔ **Robust to deformations**
- ➔ **Topological invariant**

$g = 0$



$g = 1$



© Henry Segerman

Properties that remain unaffected under smooth distortions

Chern number

$$C = \frac{1}{2\pi} \oint_{BZ} \nabla_{\mathbf{k}} \times \langle \psi(\mathbf{k}) | i \nabla_{\mathbf{k}} | \psi(\mathbf{k}) \rangle \cdot d\mathbf{s}$$

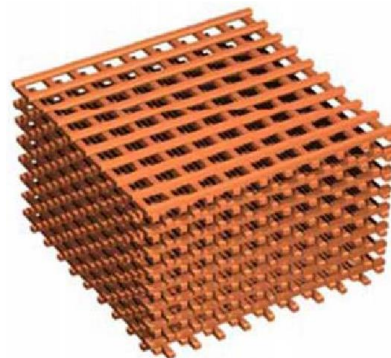
➔ $C = 0, \pm 1, \pm 2, \pm 3, \dots$

➔ Global property

➔ Robust to deformations

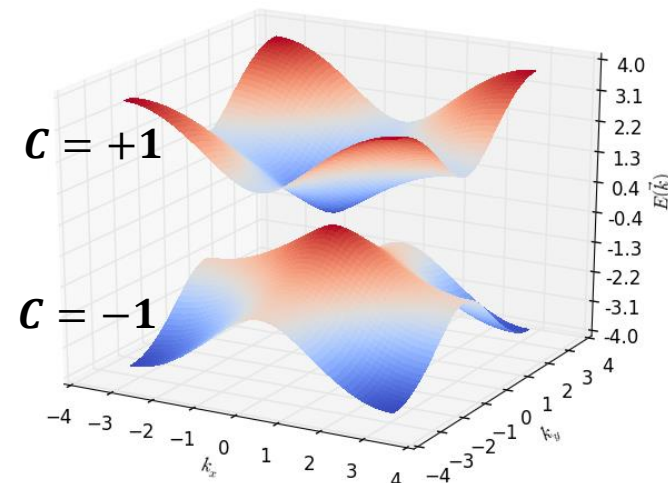
➔ Topological invariant

Electronic, phononic
or photonic crystal



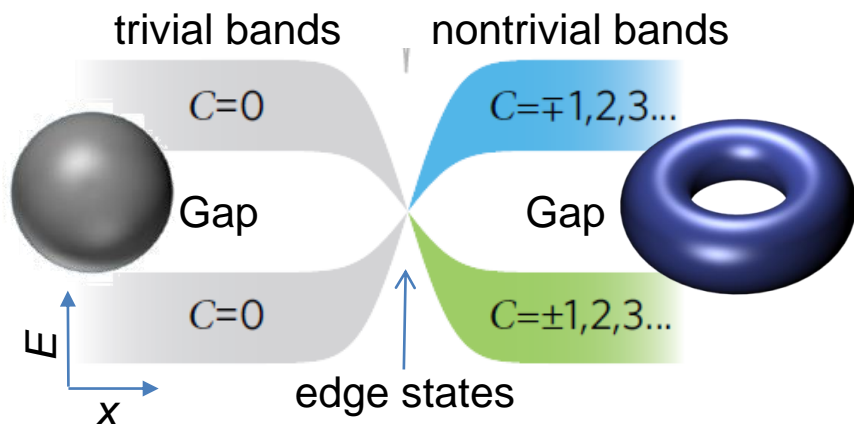
@ Intechopen

Bands



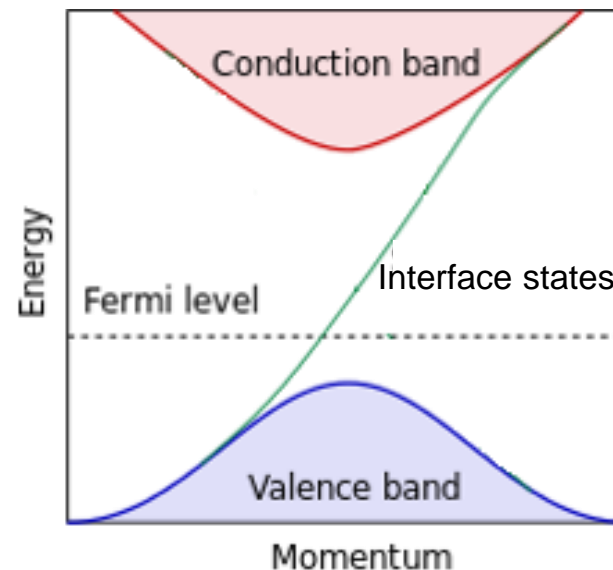
Topological invariant (Chern #)

$$C = \frac{1}{2\pi} \oint \nabla_{\mathbf{k}} \times \langle u(\mathbf{k}) | i \nabla_{\mathbf{k}} | u(\mathbf{k}) \rangle \cdot d\mathbf{s}$$



Bulk-edge correspondence

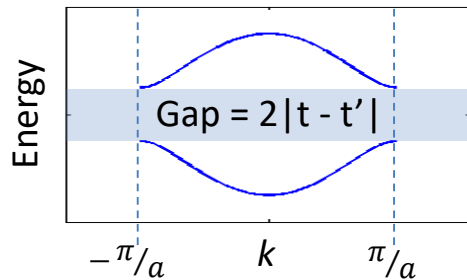
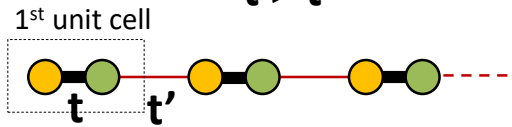
Unidirectional transport



1D lattice: the Su-Schrieffer-Heeger Hamiltonian

Dimerization 1

$$t > t'$$



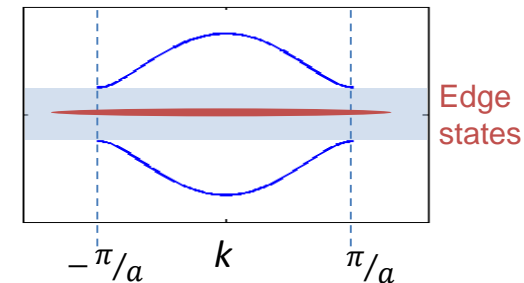
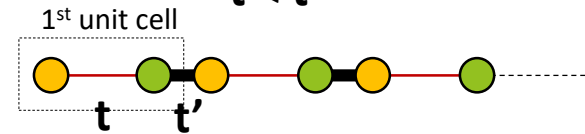
Winding number:

$$\mathcal{W} = \frac{1}{2\pi} \oint_{BZ} dk \langle \psi_{\pm}(k) | \frac{d}{dk} | \psi_{\pm}(k) \rangle = 0$$

NO edge state

Dimerization 2

$$t < t'$$



Winding number:

$$\mathcal{W} = 1$$

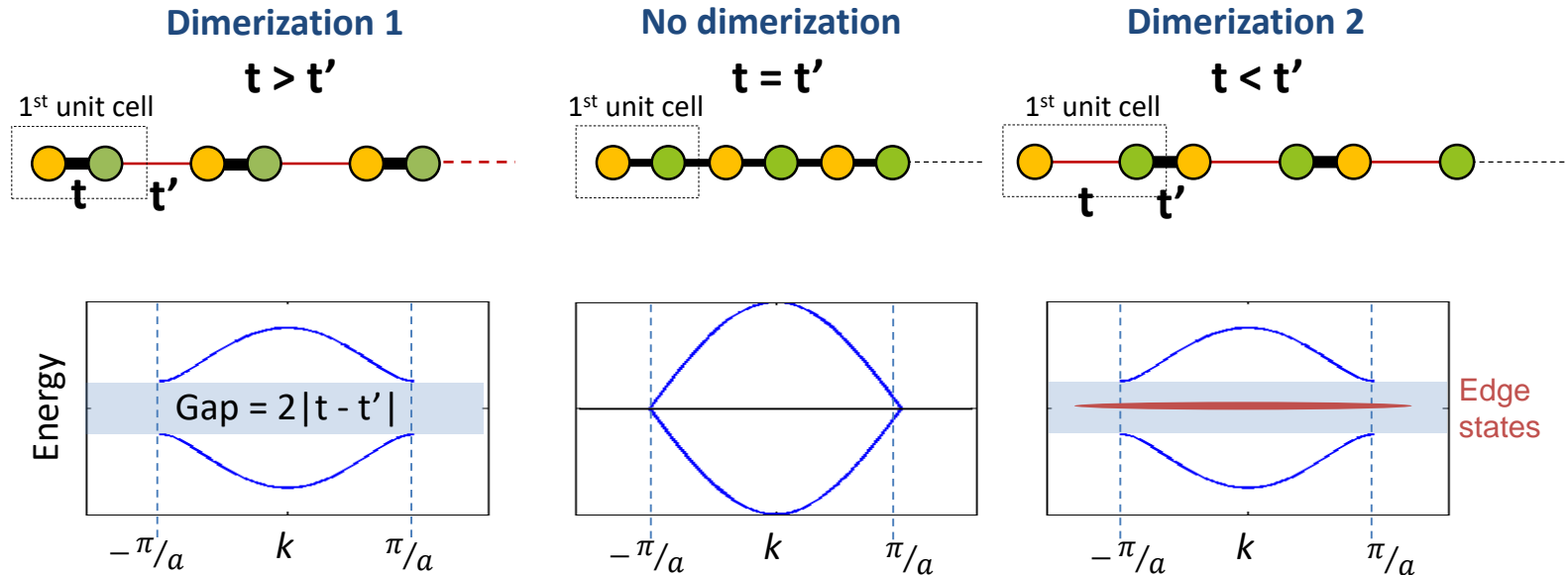
Edge states in the gap
(E=0)

Same dispersion
Eigenfunctions are different

$$|\pm\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm e^{-i\phi(k)} \end{pmatrix}$$

$$\cot\phi(k) = \frac{t'/t}{\sin ka} + \cot ka$$

1D lattice: the Su-Schrieffer-Heeger Hamiltonian



Winding number:

$$\mathcal{W} = \frac{1}{2\pi} \oint_{BZ} dk \langle \psi_{\pm}(k) | \frac{d}{dk} | \psi_{\pm}(k) \rangle = 0$$

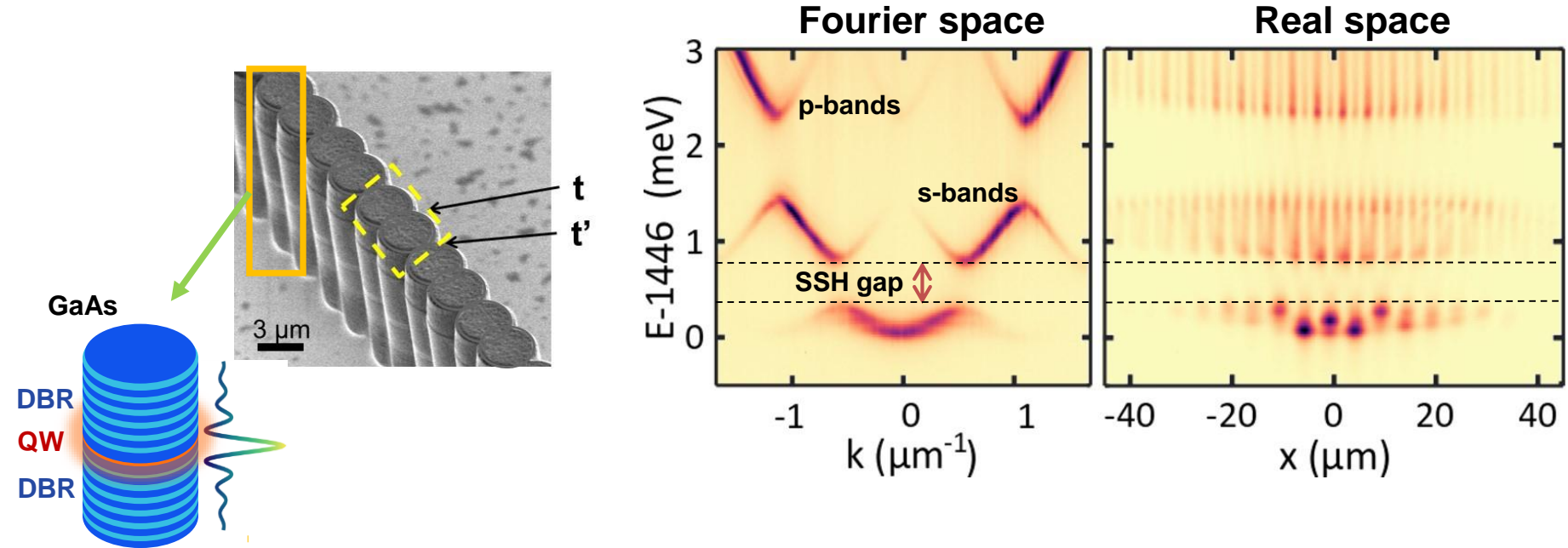
NO edge state

Winding number:

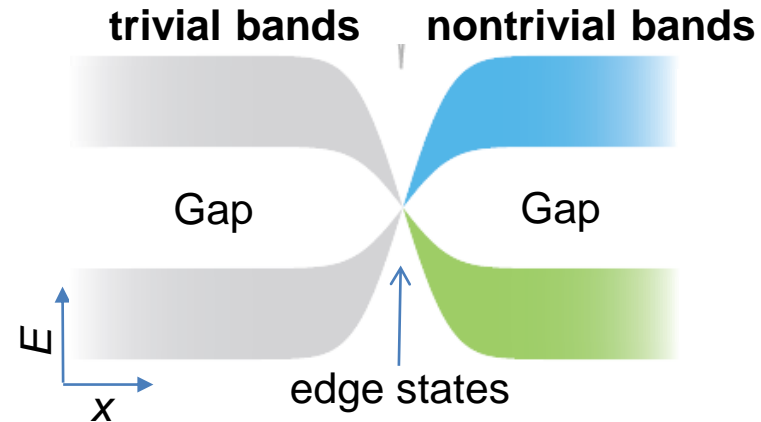
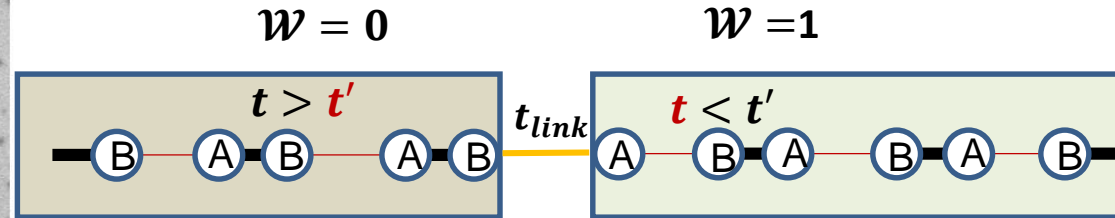
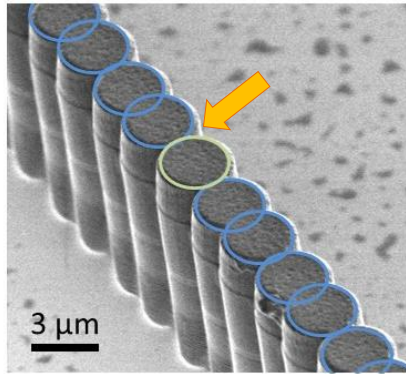
$$\mathcal{W} = 1$$

**Edge states in the gap
(E=0)**

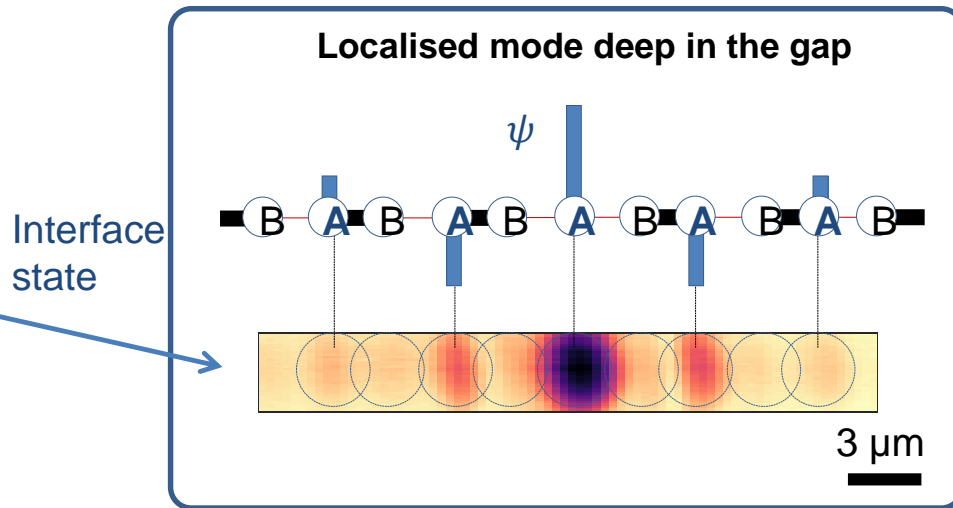
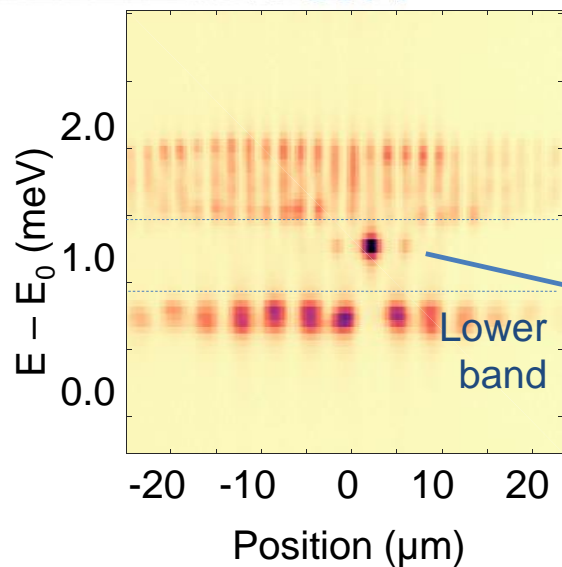
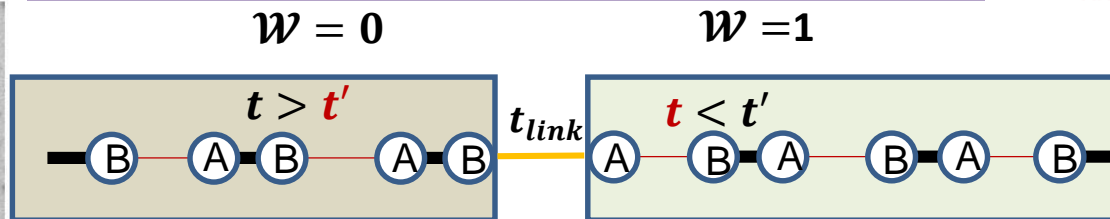
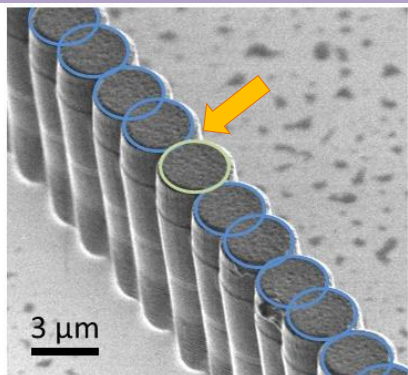
SSH lattice of photonic resonators



SSH lattice of photonic resonators



SSH lattice of photonic resonators



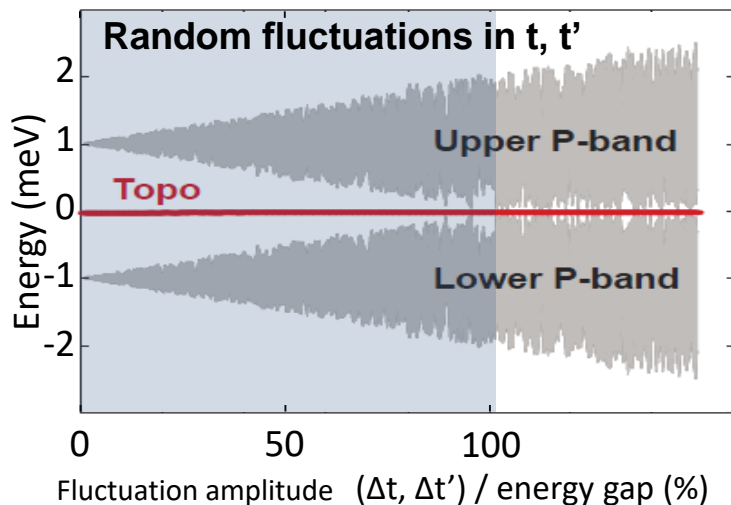
$$\{H, \sigma_z\} = 0 \text{ (chiral symmetry)}$$



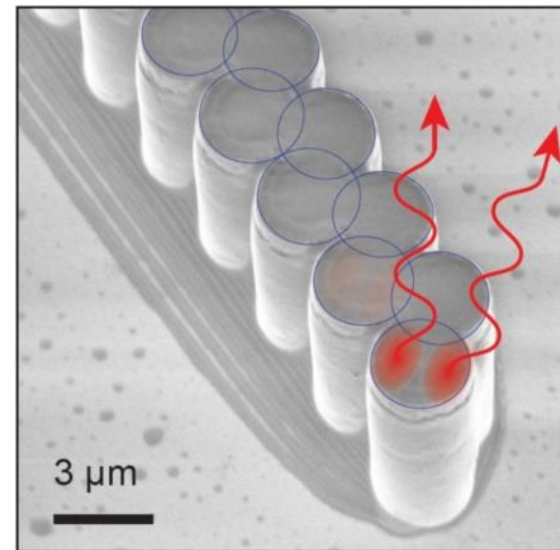
Eigspectrum is symmetric around $E=0$

Preserved even if t and t' fluctuate

Simulation



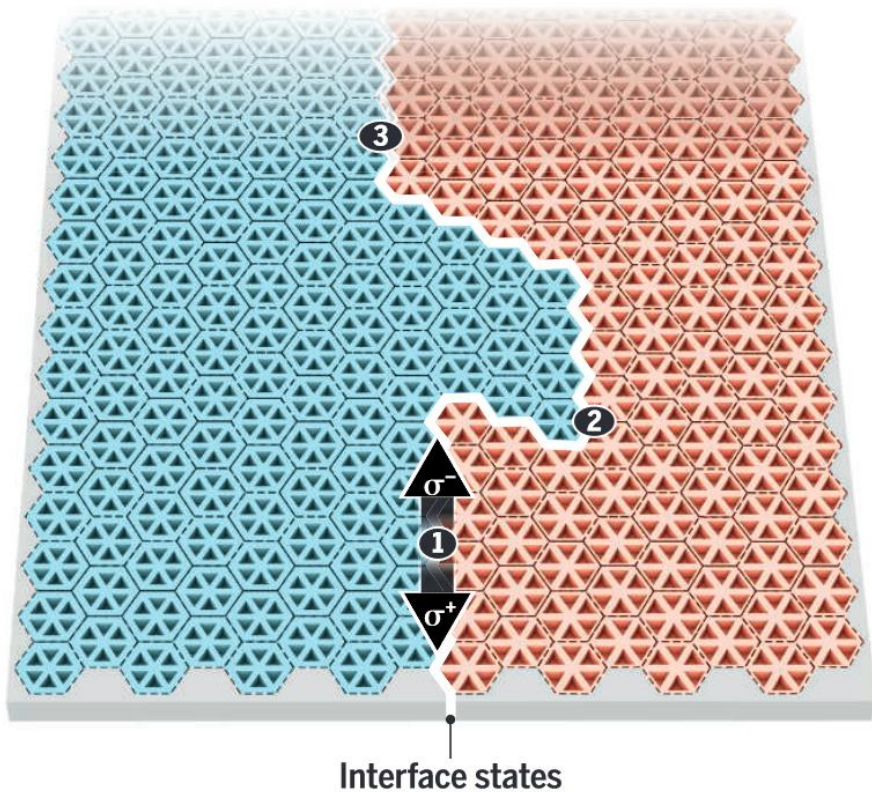
Lasing at an edge state



P. St-Jean et al., Nat. Photon. **11**, 651 (2017)

See also: H. Zhao et al., Nat. Comm. **9**, 981(2018)
 M. Parto et al., PRL **120**, 113901 (2018) } **1D**

B. Bahari et al., Science **358**, 636 (2017)
 M. A. Bandres et al., Science **359**, aar4005 (2018)
 S. Klembt et al., Nature **562**, 552 (2018) } **2D**




A. Amo, Science **359**, 638 (2018)

Is it possible to create a topological 2D material?

$$C = \frac{1}{2\pi} \oint_{BZ} \underbrace{\nabla_{\mathbf{k}} \times \langle \psi(\mathbf{k}) | i \nabla_{\mathbf{k}} | \psi(\mathbf{k}) \rangle}_{\Omega(\mathbf{k})} \cdot d\mathbf{s}$$

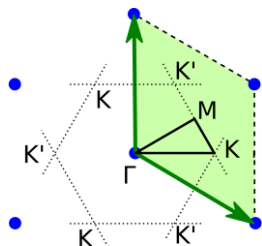
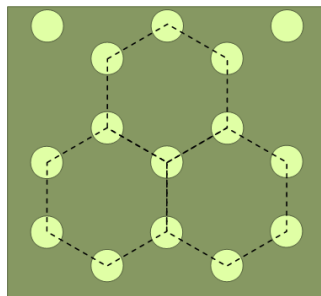
$\Omega(\mathbf{k})$ Berry curvature

- **Time-reversal symmetry** $\rightarrow \Omega(\mathbf{k}) = -\Omega(-\mathbf{k})$
 See: Z. Wang et al., Nature 461, 772 (2009)
- **Inversion symmetry** $\rightarrow \Omega(\mathbf{k}) = \Omega(-\mathbf{k})$


 $C = 0$

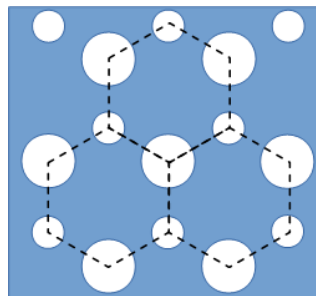
Valley Hall topology in photonic crystals

Honeycomb lattice

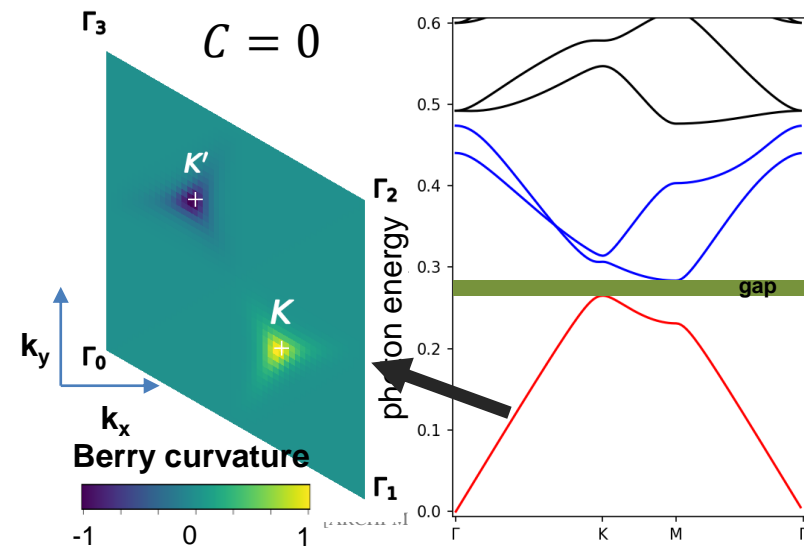
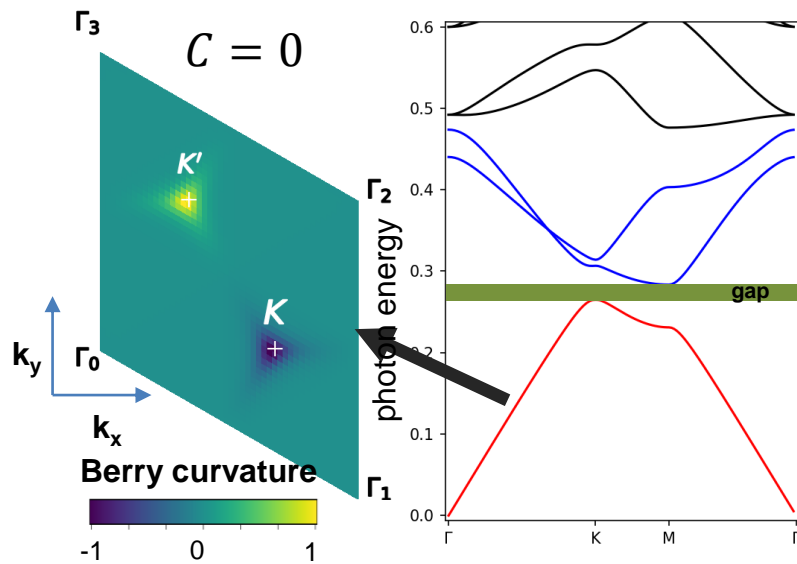
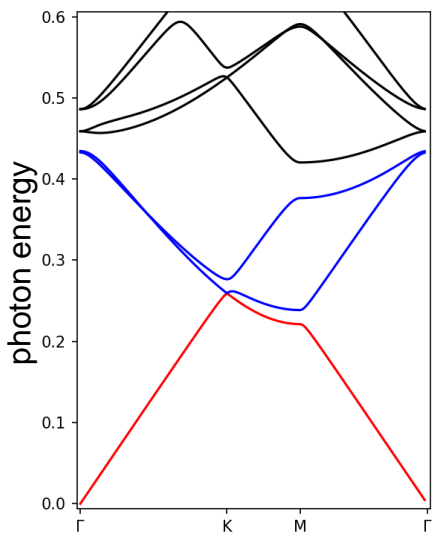
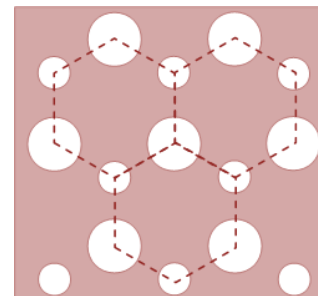


Brillouin zone

Boron nitride - 1

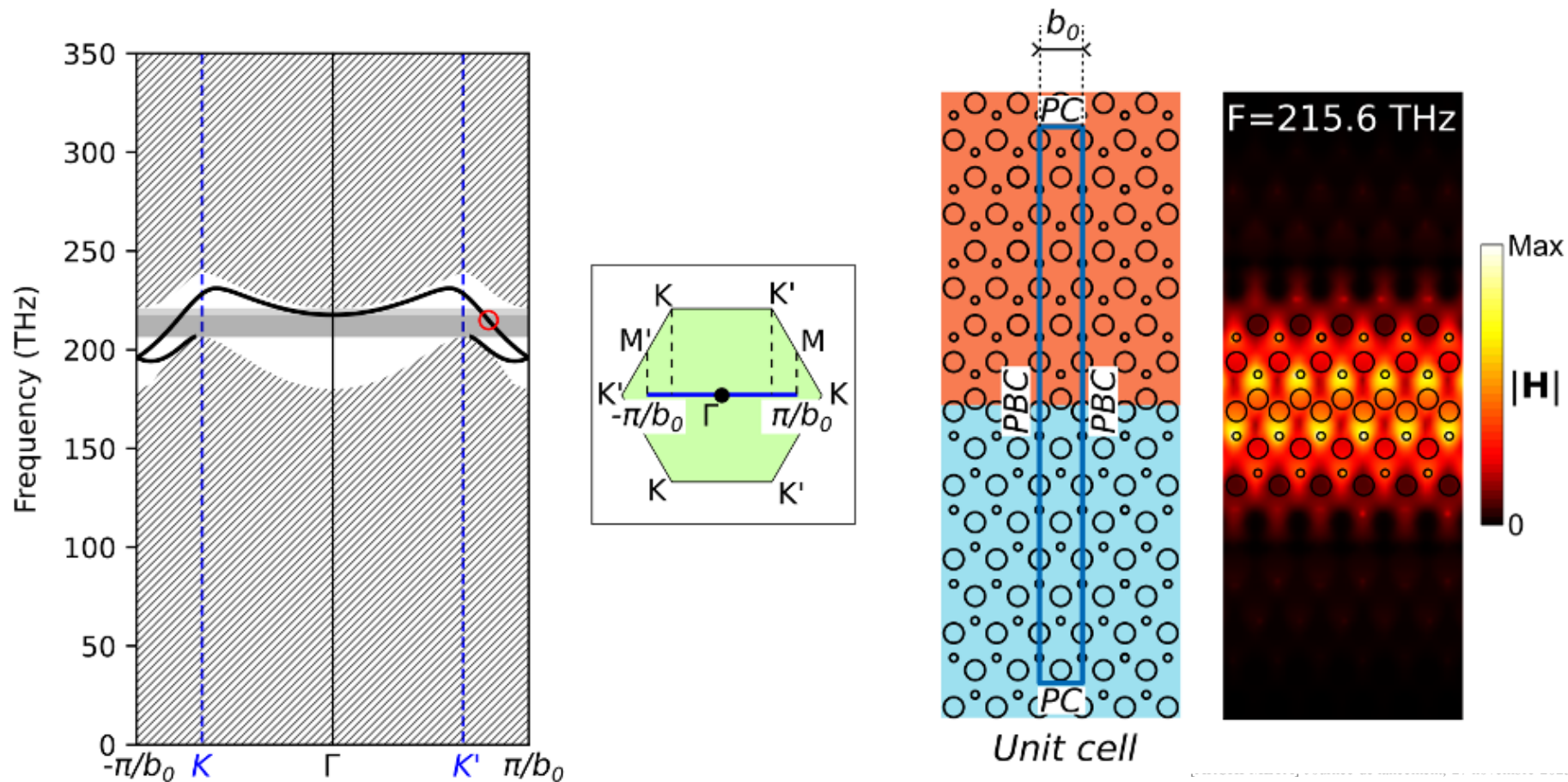


Boron nitride - 2



Valley Hall topology in photonic crystals

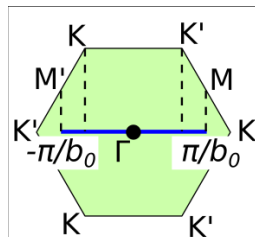
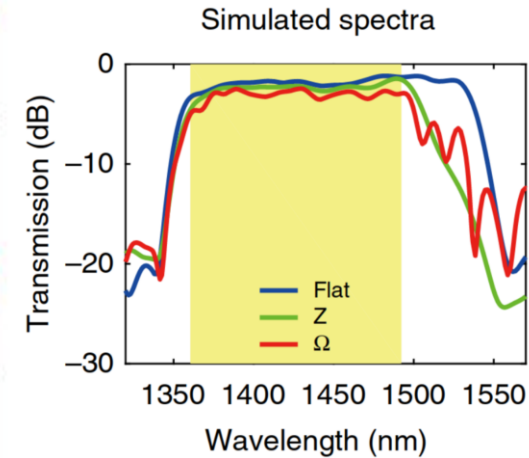
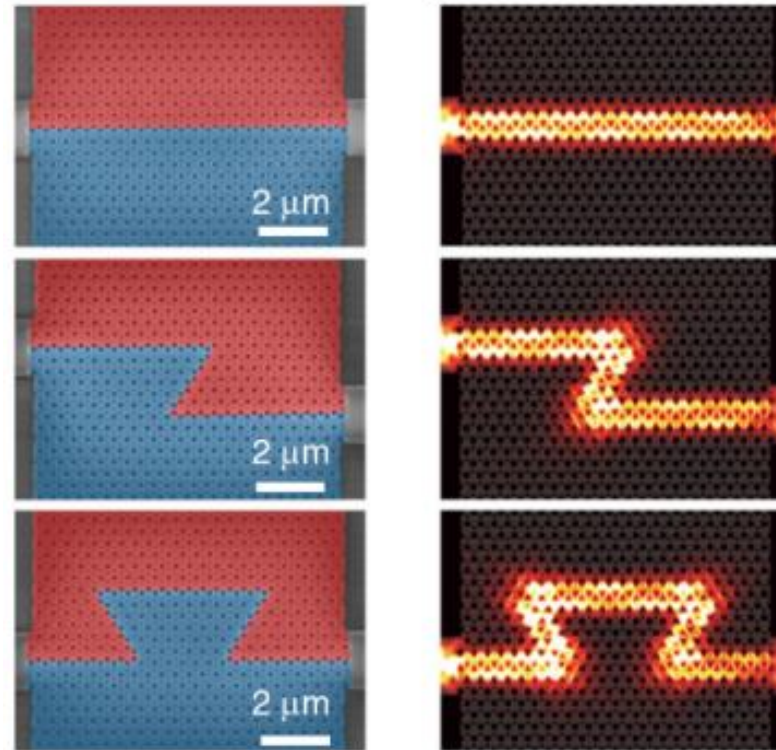
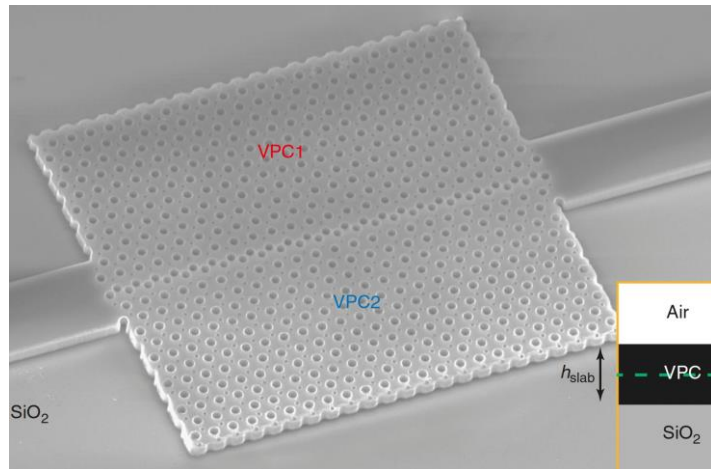
Interface modes



Valley Hall topology in photonic crystals

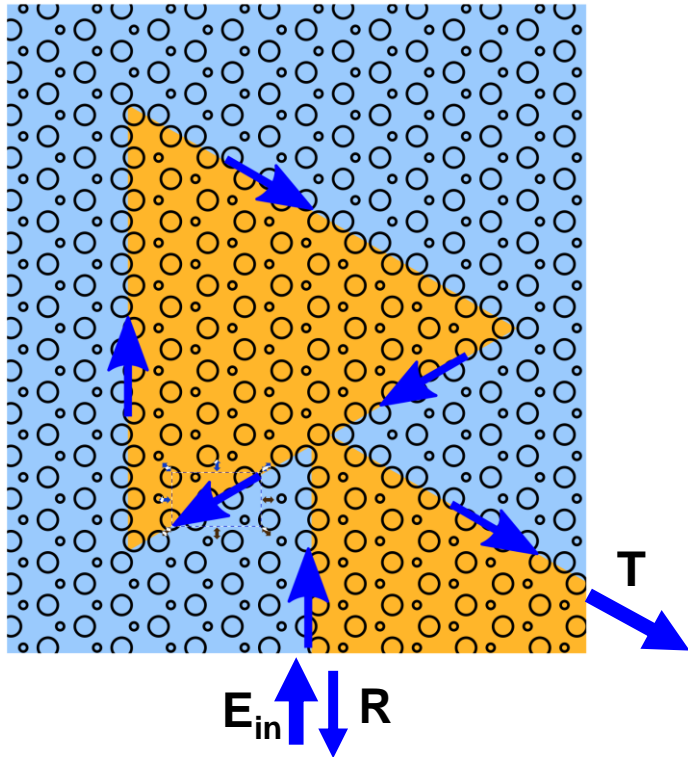
Interface modes

Go around corners with high transmission!!



As long as valley is preserved (60° turns)

Triangular cavity



If perfect topological protection

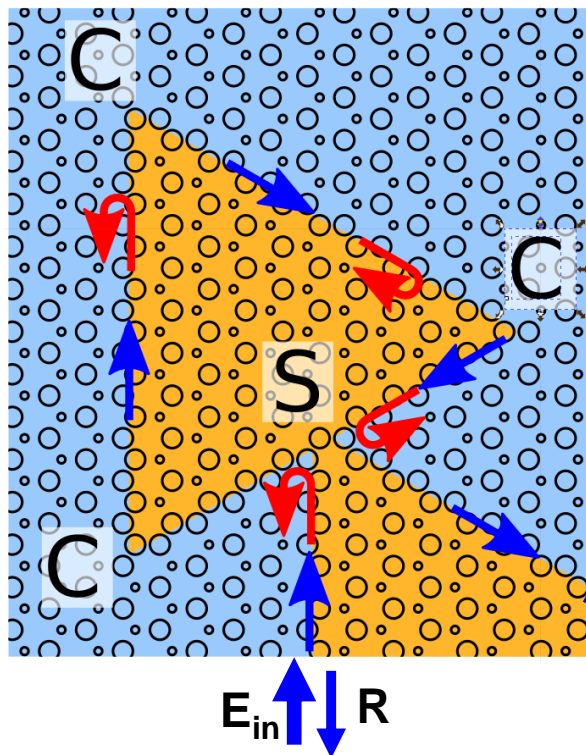


no backscattering



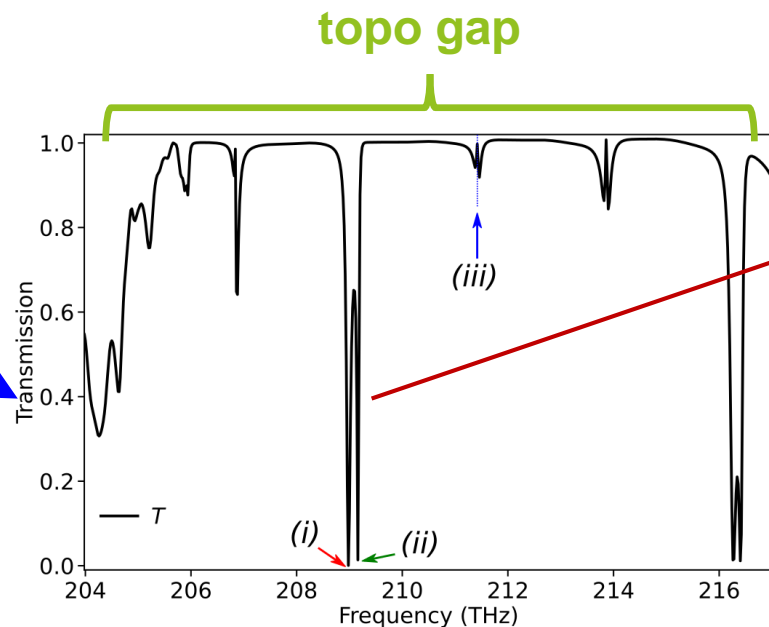
$T=1$ all over the gap

Triangular cavity

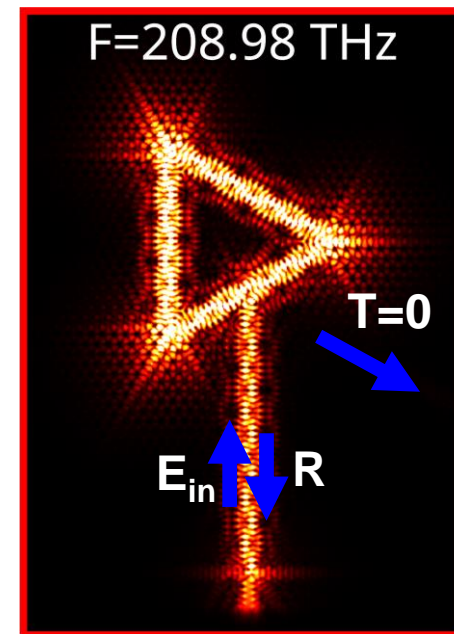


From finite element simulations:
backreflection at corners $\sim 1-2\%$

Triangular holes: 0.01%



simulation



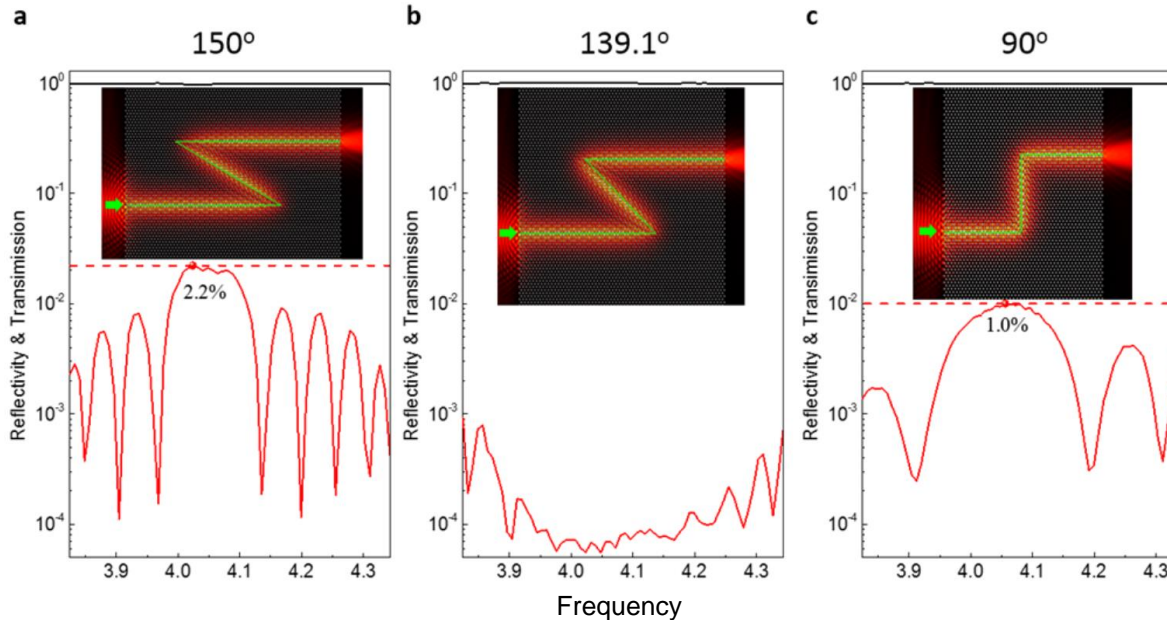
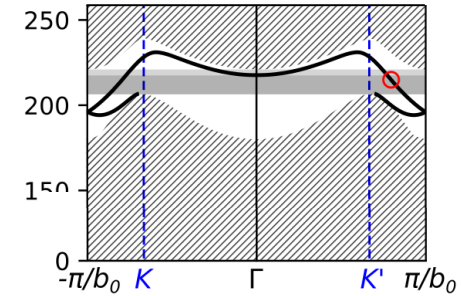
Topological protection is not perfect

G. Lévêque et al., PRA **108**, 043505 (2023)

See also: S. Arora et al., LSA **10**, 9 (2021)

C. A. Rosiek et al., Nat. Photon. **17**, 386 (2023)

K-valley conservation is not required!!

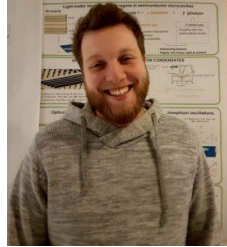


Is topology the right framework to understand these channels?



Palaiseau
(France)

Lattices of micropillars



Nicolas
Pernet



Philippe
St-Jean



Nicola
Carlon Zambon



Jacqueline
Bloch



Sylvain
Ravets



Martina
Morassi



Aristide
Lemaître



Luc
LeGratiet



Isabelle
Sagnes



Abdelmounaim
Harouri

Sample fabrication

Photonic crystals



G. Levecque
IEMN



Y. Pennec
IEMN



P. Sznitgiser
PhLAM



A. Martinez
UP Valencia